

# Beyond Closure Models: Learning Chaotic Systems via Physics-Informed Neural Operators

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## Background

**Central Task:** Estimate Long-term Statistics of Chaotic Systems with Coarse-grid Simulations

$$\begin{cases} \partial_t u(x, t) = \mathcal{A}u(x, t) \\ u(x, 0) = u_0(x), u_0 \in \mathcal{H} \end{cases}$$

$\mathcal{A}$ : (Nonlinear) Operator;  $\mathcal{H}$ : function space of interest.

Attractor  $\Omega$ : All trajectories  $\{u(\cdot, t)\}$  will converge to it as  $t \rightarrow \infty$ .

Invariant Measure:  $\mu^* := \lim_{T \rightarrow \infty} \frac{1}{T} \int_{t=0}^T \delta_{S(t)u} dt, u \in \mathcal{H}, a.e.$  (average along traj).

**Long-term Statistics:**  $\mathbb{E}_{u \sim \mu^*} \mathcal{O}(u)$  for measurement functionals  $\mathcal{O}$ .

- Key information of the physical system at dynamical equilibrium.
- Important in application: airfoil design, climate modeling, etc.
- [REMARK] Impossible to track trajectories for very-long time in chaotic system, but possible to estimate statistics (shadowing lemma).

### General Methods:

- Straightforward: Fully-Resolved Simulations (FRS)  
Numerically simulate with very fine spatio-temporal grids/meshes.  
**TOO EXPENSIVE!** Intractable for most practical problems.
- Estimations statistics with coarse-grid simulation:  
Need to account for the large discretization error (i.e. missing information from the fine scale).
- Known as **Closure Modeling** or **Coarse-graining**.

### Scheme of Closure Modeling:

- $\mathcal{F}$ : filter from fine grid to coarse grid, e.g. spatial downsampling, Fourier mode truncation etc.), viewed as a mapping in function space  $\mathcal{H}$ .
- Filtered Dynamic  $\partial_t \bar{u} = \mathcal{F}\mathcal{A}u = \mathcal{A}\bar{u} + \underbrace{(\mathcal{F}\mathcal{A} - \mathcal{A}\mathcal{F})u}_{\text{Unresolved}}$ , ( $\bar{u} := \mathcal{F}u$ ).

$$\begin{cases} \partial_t v(x, t) = \mathcal{A}v(x, t) + \text{clos}(v; \theta), x \in D' \\ v(x, 0) = \bar{u}_0(x), \bar{u}_0 \in \mathcal{F}(\mathcal{H}), \end{cases}$$

- Interpretation:
- Assign a vector field ( $\mathcal{A} + \text{clos}$ ) in the reduced space to drive the dynamics.

### How to design closure models?

- Classical Models: hand-designed. Strong physical intuition and assumptions.
- Machine Learning for Closure Models (Hopes: better expressiveness)
- [Learning Framework] Supervised Learning (Single-State Model)

$$J_{ap}(\theta; \mathcal{D}) = \frac{1}{|\mathcal{D}|} \sum_{i \in \mathcal{D}} \|\text{clos}(\bar{u}_i; \theta) - (\mathcal{F}\mathcal{A} - \mathcal{A}\mathcal{F})u_i\|^2$$

$u_i$ : data from fully-resolved (fine-grid) simulations

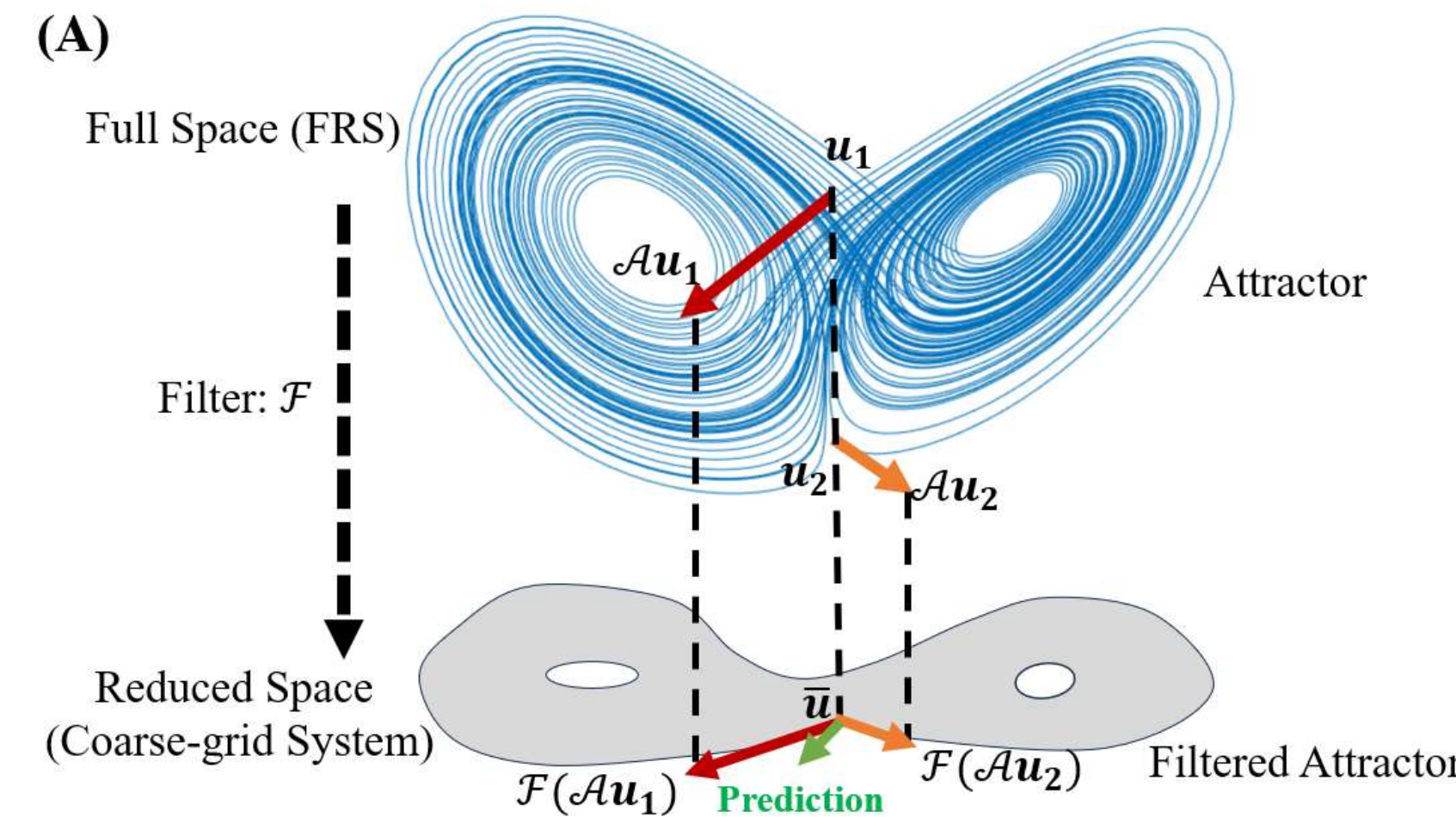
### [Advanced Variants]

- [Posterior Training]  $J_{post}(\theta; \mathcal{D}) = J_{ap}(\theta) + \frac{1}{|\mathcal{D}|} \sum_{i \in \mathcal{D}} \|v_i(\cdot, \Delta t; \theta) - \mathcal{F}(S(\Delta t)u_i)\|^2$

- [History-aware Models] Model's input:  $\{\bar{u}(x_i, t-s)\}_{x_i \in D', 0 < s \leq t_0}$
- [Stochastic Closure Models]

## Theoretical Results

**Learning-based methods should not follow previous closure modeling ansatz  $\mathcal{A}\bar{u} + \text{clos}(\bar{u}; \theta)$ .**



- Learning-based closure models suffers from a **large approximation error** independent of model complexity, stemming from the **non-uniqueness** of the target mapping.
- Leveraging history information and randomness can neither help.
- A fundamental limitation for any method following the ansatz  $\mathcal{A} + \text{clos}_\theta$ .
- To mitigate the nonunique issue, **model has to use a large number** of FRS data. The amount of training data is of the same order to estimate long-term statistics! – **eliminating the need for a closure model!**
- One could not expect the model to generalize among different dynamics (e.g. different domain shape, different coefficient in the PDE).

### Proof Idea: Functional Liouville Flow

- View functions  $u$  as particles.
- (Infinite-dimensional) Liouville eqn. for analyzing the limit distribution.

Coarse-grid Dynamics that can achieve optimal approx. of  $\mu^*$ :

$$\partial_t v = \mathbb{E}_{u \sim \mu_t} [\mathcal{F}\mathcal{A}u | \mathcal{F}u = v], \mu_t: \text{distribution of } u \in \mathcal{H} \text{ at time } t.$$

For CGS and learning closure model, one can only fix a  $\hat{\mu} \in \mathcal{P}(\mathcal{H})$ , and evolve  $\partial_t v = \mathbb{E}_{u \sim \hat{\mu}} [\mathcal{F}\mathcal{A}u | \mathcal{F}u = v]$

- $\hat{\mu} = \mu^*$ : optimal approximation of  $\mu^*$  in reduced system ( $\mathcal{F}\# \mu^*$ ).
- In practice, the best model one can yield (assuming sufficient expressive power of NN function class) corresponds to  $\hat{\mu} = \mu_{data}$ .
- Large gap between  $\mu_{data}$  and  $\mu^*$  for infinite-dim distributions!

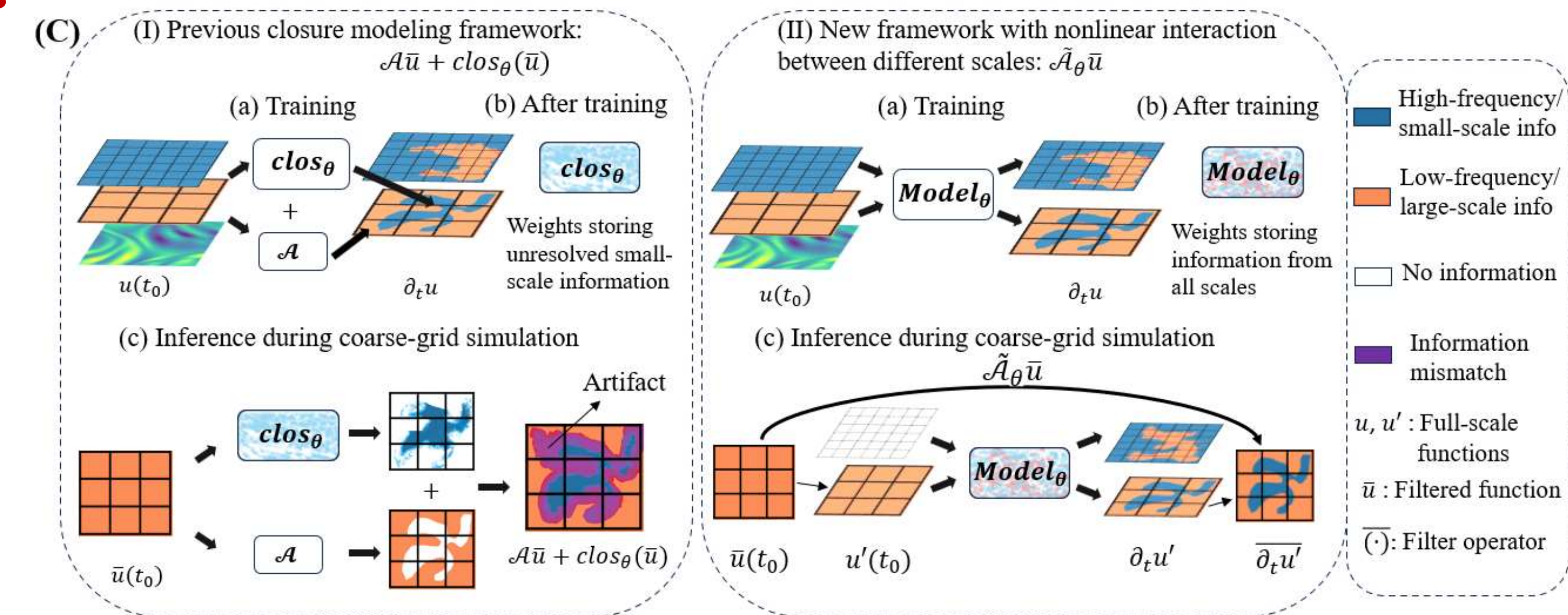
### Key Takeaway

➢ We need **nonlinear** interaction between information from different scales (i.e. resolved part in coarse-grid system and unresolved parts!)

➢ Previous ansatz:  $\mathcal{A} + \text{clos}(\cdot, \theta)$

➢ New Ansatz:  $\tilde{\mathcal{A}}_\theta$

## New Ansatz with Physics-Informed Operator Learning



### Neural Operators

- **Resolution-invariant.** (Support input from both coarse-grid and fine-grid).
- **$O(1)$  jump along time** instead of moving with tiny time grids.
- The infinitesimal generator of learned operator  $\mathcal{G}_\theta$  plays the role of  $\tilde{\mathcal{A}}_\theta$
- Physics-informed learning + multi-resolution pre-training to reduce reliance on FRS data (**only  $\sim 10^2$  snapshots from single FRS trajectory** vs  $\sim 10^5$  in previous works).
- Theoretical guarantee on optimal estimation of  $\mu^*$  with coarse-grid simulations.

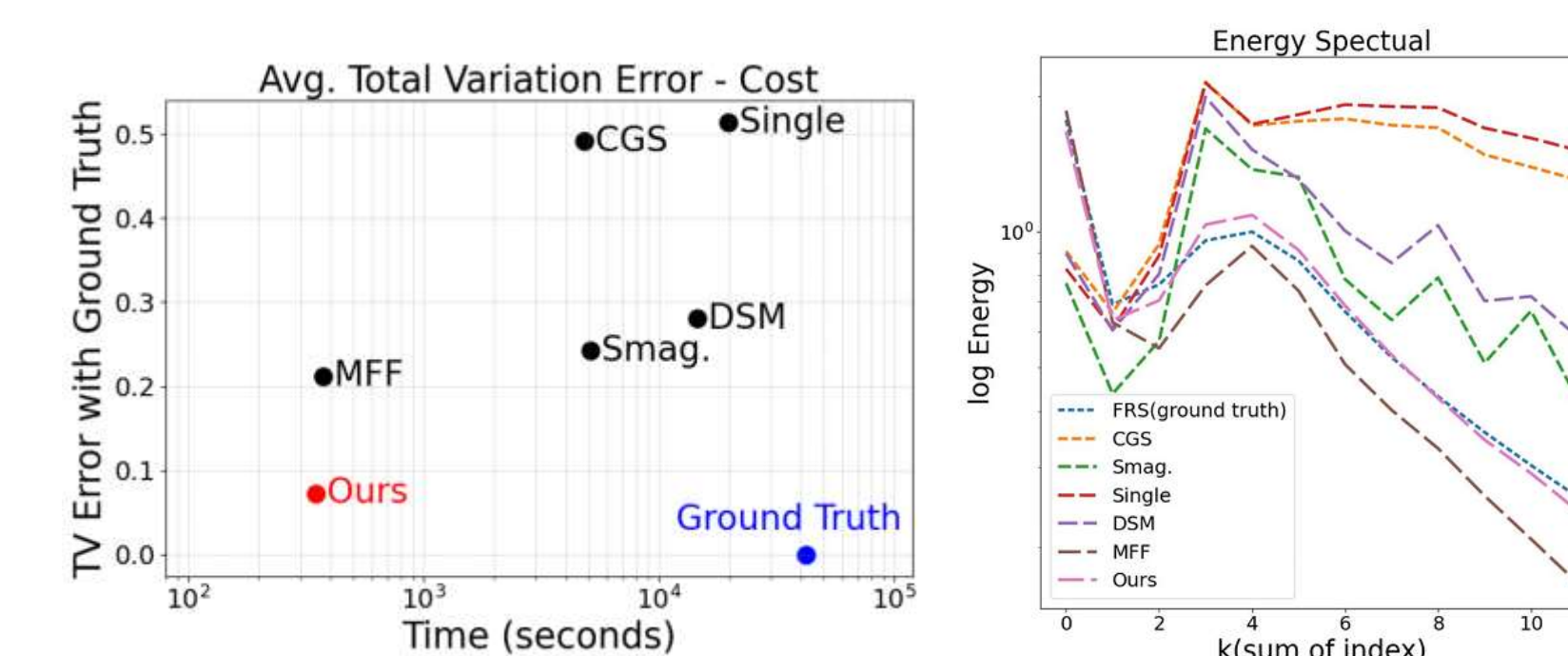
**Theorem 3.1.** For any  $h > 0$ , denote  $\hat{\mu}_{h,\theta} := \lim_{N \rightarrow \infty} \frac{1}{N} \sum_{n=1}^N \delta_{\mathcal{G}_\theta^n v_0(x)}$ , any  $v_0(x)$  with  $x \in D'$ . For any  $\epsilon > 0$ , there exists  $\delta > 0$  s.t. **as long as  $\|(\mathcal{G}_\theta u)(\cdot, h) - S(h)u\|_{\mathcal{H}} < \delta, \forall u \in \mathcal{H}$ , we have  $\mathcal{W}_{\mathcal{H}}(\hat{\mu}_{h,\theta}, \mathcal{F}\# \mu^*) < \epsilon$** , where  $\mathcal{W}_{\mathcal{H}}$  is a generalization of Wasserstein distance in function space.

## (Part of) Experiment Results

Kolmogorov Flow (2D forced Navier-Stokes eqn).

$$\partial_t \mathbf{u} = -(\mathbf{u} \cdot \nabla) \mathbf{u} - \nabla p + \nu \Delta \mathbf{u} + (\sin(4y), 0)^T, \quad \nabla \cdot \mathbf{u} = 0, \quad (x, y, t) \in [0, L]^2 \times \mathbb{R}_+$$

Method	Avg. TV	Energy	Vorticity	Variance	FRS	CGS (No closure)
CGS (No closure)	0.4914	178.4651%	0.1512	253.4234%	39.70	4.50
Smagorinsky	0.2423	52.9511%	0.0483	20.1740%	4.81	4.81
Single-state	0.5137	205.3709%	0.1648	298.2027%	18.57	18.57
DSM	0.2803	74.2150%	0.0821	73.6158%	13.67	13.67
MFF	0.2123	20.7055%	0.0115	20.4410%	<b>0.32</b>	<b>0.32</b>
<b>Our Method</b>	<b>0.0726</b>	<b>5.3276%</b>	<b>0.0091</b>	<b>2.8666%</b>	<b>0.32</b>	<b>0.32</b>



The newest version:  
[arxiv.org/abs/2408.05177](https://arxiv.org/abs/2408.05177)